



Two-Dimensional Phase-Field Simulation of Self-Assembled Quantum Dot Formation

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INTRODUCTION

Phase-field simulations are performed to investigate the effects of the growth parameters on island formation and island morphological evolution during deposition for SiGe/Si(001) system. In the phase-field model, facet morphologies are modeled using the generalized gradient correction coefficient for a crystal with a high anisotropy of surface energy. Furthermore, to enable computation of the island morphological change from faceted pyramids to multifaceted domes with increasing Ge coverage, we apply the sixteen-fold of the surface energy. The validity of the model developed is confirmed by comparing the numerical results demonstrated here with the reported experimental observations.

PHASE-FIELD MODEL

< Free-energy functional >

The morphologies of the thin film are assumed to be determined by the competition of the surface energy and the elastic strain energy. The Ginzburg-Landau-type free-energy functional of the system is assumed to have the form

$$F = \int_V \left\{ f_d(\phi) + f_e(\phi, \varepsilon_{ij}) + \frac{a^2}{2} |\nabla \phi|^2 \right\} dV$$

where

$$\phi : \text{phase field} \quad \phi = \begin{cases} 0 : \text{vapor phase} \\ 1 : \text{solid phase (= substrate + film)} \end{cases}$$

f_d : double-well potential

$$f_d(\phi) = Wg(\phi) \quad g(\phi) = \phi^2(1-\phi)^2$$

f_e : elastic strain energy

$$f_e(\phi, \varepsilon_{ij}) = \frac{1}{2} D_{ijkl}(\phi) (\varepsilon_{ij} - \varepsilon_{ij}^0) (\varepsilon_{kl} - \varepsilon_{kl}^0)$$

$D_{ijkl}(\phi) = \rho(\phi) D_{ijkl}^0$: elastic coefficient

$$\rho(\phi) = \frac{1}{2} \left[\tanh \frac{2\phi-1}{2\tau} + 1 \right] : \text{density of solid phase}$$

a : gradient correction coefficient

< Faceted island >

The surface anisotropy is considered by

$$a(\theta) = \bar{a} \{1 + \gamma \cos k\theta\}$$

where

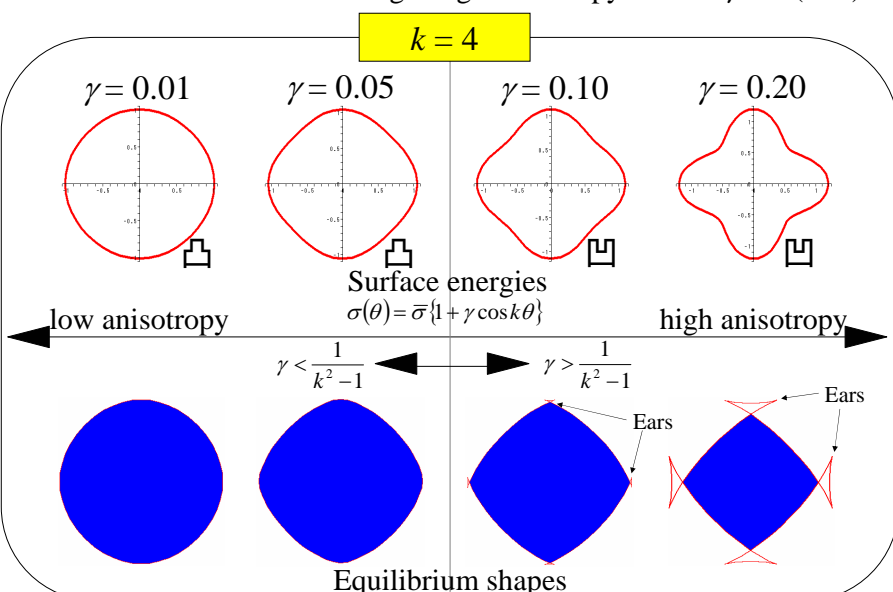
\bar{a} : constant related to the surface energy σ and the surface thickness δ

γ : strength of anisotropy

k : mode number

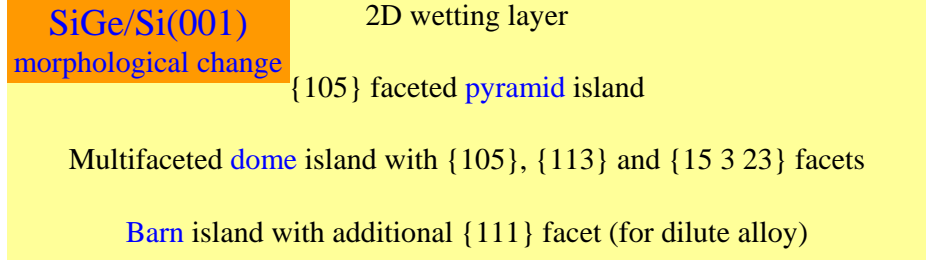
θ : angle between the surface normal and the x -axis

Faceted islands are created using a high anisotropy such as $\gamma > 1/(k^2-1)$

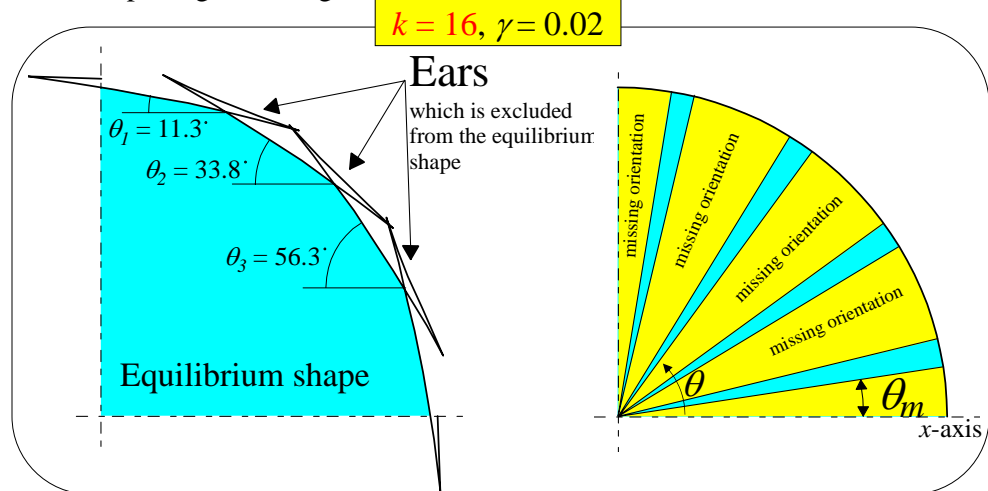


< Multifaceted island >

SiGe/Si(001) shows the following surface morphological change:



We introduce $k = 16$ to simulate the multifaceted islands and island morphological change.



The first missing orientation θ_m is calculated from

$\theta_1 = 11.3^\circ$	{105} facet plane
$\theta_2 = 33.8^\circ$	{15 3 23} facet plane
$\theta_3 = 56.3^\circ$	{111} facet plane

$$a(\theta_m) \sin \theta_m + a_\theta(\theta_m) \cos \theta_m = 0$$

< Generalized gradient coefficient >

To provide the island morphology without "ears" corresponding to the missing orientation, we use the regularized gradient coefficient for the missing orientation.

$$a(\theta) = \begin{cases} \bar{a} \{1 + \gamma \cos k\theta\} & \text{for } (2\pi i/k + \theta_m) \leq \theta \leq (2\pi(i+1)/k - \theta_m) \\ \frac{a(\theta_m)}{\cos \theta} \cos \theta & \text{for } (2\pi i/k - \theta_m) < \theta < (2\pi i/k + \theta_m) \\ & \text{(missing orientation)} \end{cases}$$

where i denotes the integers from 0 to $k-1$.

< TDGL (Time-dependent Ginzburg-Landau) Equation >

$$\frac{\partial \phi}{\partial t} = \nabla \cdot \left(M \nabla \frac{\delta F}{\delta \phi} \right) + V_d n_y \chi$$

$$= M \nabla^2 \left[2\phi(1-\phi)(1-2\phi)W + \frac{\partial f_e(\phi, \varepsilon_{ij})}{\partial \phi} + \frac{\partial}{\partial x} \left(a \frac{\partial a}{\partial \theta} \frac{\partial \phi}{\partial y} \right) - \frac{\partial}{\partial y} \left(a \frac{\partial a}{\partial \theta} \frac{\partial \phi}{\partial x} \right) - a^2 \nabla^2 \phi \right] + V_d n_y \chi$$

$$\sigma_{ij,j} = 0 : \text{stress equilibrium condition}$$

where

M : mobility representing the surface diffusion

V_d : deposition rate

n_y : y -direction component of the surface normal

χ : random number

The phase field ϕ is assumed to be a conserved parameter. We assume the stress equilibrium conditions $\sigma_{ij} = 0$, since the elastic relaxation occurs much faster than the surface diffusion.

